

Suppose the rectangular region in Figure 12.23(a) is a thin plate in which the temperature  $u$  is a function of time  $t$  and position  $(x, y)$ . Then, under suitable conditions,  $u(x, y, t)$  can be shown to satisfy the **two-dimensional heat equation**

$$k \left( \frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} \right) = \frac{\partial u}{\partial t}. \quad (1)$$

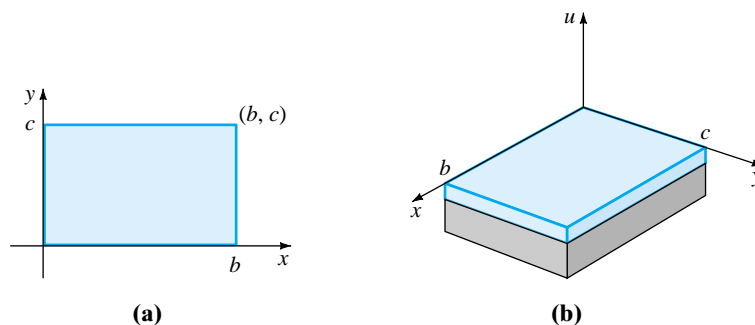
On the other hand, suppose Figure 12.23(b) represents a rectangular frame over which a thin flexible membrane has been stretched (a rectangular drum). If the membrane is set in motion, then its displacement  $u$ , measured from the  $xy$ -plane (transverse vibrations), is also a function of  $t$  and position  $(x, y)$ . When the vibrations are small, free, and undamped,  $u(x, y, t)$  satisfies the **two-dimensional wave equation**

$$a^2 \left( \frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} \right) = \frac{\partial^2 u}{\partial t^2}. \quad (2)$$

To separate variables in (1) and (2), we assume a product solution of the form  $u(x, y, t) = X(x)Y(y)T(t)$ . We note that

$$\frac{\partial^2 u}{\partial x^2} = X''YT, \quad \frac{\partial^2 u}{\partial y^2} = XY''T, \quad \text{and} \quad \frac{\partial u}{\partial t} = XYT'.$$

As we see next, with appropriate boundary conditions, boundary-value problems involving (1) and (2) lead to the concept of Fourier series in two variables.



**FIGURE 12.23** (a) Rectangular plate and (b) rectangular membrane

Find the temperature  $u(x, y, t)$  in the plate shown in Figure 12.23(a) if the initial temperature is  $f(x, y)$  throughout and if the boundaries are held at temperature zero for time  $t > 0$ .

**SOLUTION** We must solve

$$k \left( \frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} \right) = \frac{\partial u}{\partial t}, \quad 0 < x < b, \quad 0 < y < c, \quad t > 0$$

$$\text{subject to} \quad u(0, y, t) = 0, \quad u(b, y, t) = 0, \quad 0 < y < c, \quad t > 0$$

$$u(x, 0, t) = 0, \quad u(x, c, t) = 0, \quad 0 < x < b, \quad t > 0$$

$$u(x, y, 0) = f(x, y), \quad 0 < x < b, \quad 0 < y < c.$$

Substituting  $u(x, y, t) = X(x)Y(y)T(t)$ , we get

$$k(X''YT + XY''T) = XYT' \quad \text{or} \quad \frac{X''}{X} = -\frac{Y''}{Y} + \frac{T'}{kT}. \quad (3)$$

Since the left-hand side of the last equation in (3) depends only on  $x$  and the right side depends only on  $y$  and  $t$ , we must have both sides equal to a constant  $-\lambda$ :

$$\frac{X''}{X} = -\frac{Y''}{Y} + \frac{T'}{kT} = -\lambda$$

and so 
$$X'' + \lambda X = 0 \quad (4)$$

$$\frac{Y''}{Y} = \frac{T'}{kT} + \lambda. \quad (5)$$

By the same reasoning, if we introduce another separation constant  $-\mu$  in (5), then

$$\frac{Y''}{Y} = -\mu \quad \text{and} \quad \frac{T'}{kT} + \lambda = -\mu$$

yield 
$$Y'' + \mu Y = 0 \quad \text{and} \quad T' + k(\lambda + \mu)T = 0. \quad (6)$$

Now the homogeneous boundary conditions

$$\left. \begin{array}{l} u(0, y, t) = 0, \quad u(b, y, t) = 0 \\ u(x, 0, t) = 0, \quad u(x, c, t) = 0 \end{array} \right\} \text{ imply that } \begin{cases} X(0) = 0, & X(b) = 0 \\ Y(0) = 0, & Y(c) = 0. \end{cases}$$

Thus we have two Sturm-Liouville problems

$$X'' + \lambda X = 0, \quad X(0) = 0, \quad X(b) = 0 \quad (7)$$

and 
$$Y'' + \mu Y = 0, \quad Y(0) = 0, \quad Y(c) = 0. \quad (8)$$

The usual consideration of cases ( $\lambda = 0$ ,  $\lambda = \alpha^2 > 0$ ,  $\lambda = -\alpha^2 < 0$ ,  $\mu = 0$ , and so on) leads to two independent sets of eigenvalues

$$\lambda_m = \frac{m^2 \pi^2}{b^2} \quad \text{and} \quad \mu_n = \frac{n^2 \pi^2}{c^2}.$$

The corresponding eigenfunctions are

$$X(x) = c_2 \sin \frac{m\pi}{b} x, \quad m = 1, 2, 3, \dots, \quad \text{and} \quad Y(y) = c_4 \sin \frac{n\pi}{c} y, \quad n = 1, 2, 3, \dots \quad (9)$$

After substituting the known values of  $\lambda_n$  and  $\mu_n$  in the first-order DE in (6), its general solution is found to be  $T(t) = c_5 e^{-k[(m\pi/b)^2 + (n\pi/c)^2]t}$ . A product solution of the two-dimensional heat equation that satisfies the four homogeneous boundary conditions is then

$$u_{mn}(x, y, t) = A_{mn} e^{-k[(m\pi/b)^2 + (n\pi/c)^2]t} \sin \frac{m\pi}{b} x \sin \frac{n\pi}{c} y,$$

where  $A_{mn}$  is an arbitrary constant. Because we have two sets of eigenvalues, we are prompted to try the superposition principle in the form of a double sum

$$u(x, y, t) = \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} A_{mn} e^{-k[(m\pi/b)^2 + (n\pi/c)^2]t} \sin \frac{m\pi}{b} x \sin \frac{n\pi}{c} y. \quad (10)$$

At  $t = 0$  we must have

$$u(x, y, 0) = f(x, y) = \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} A_{mn} \sin \frac{m\pi}{b} x \sin \frac{n\pi}{c} y. \quad (11)$$

We can find the coefficients  $A_{mn}$  by multiplying the double sum (11) by the product  $\sin(m\pi x/b) \sin(n\pi y/c)$  and integrating over the rectangle  $0 \leq x \leq b$ ,  $0 \leq y \leq c$ . It follows that

$$A_{mn} = \frac{4}{bc} \int_0^c \int_0^b f(x, y) \sin \frac{m\pi}{b} x \sin \frac{n\pi}{c} y \, dx \, dy. \quad (12)$$

Thus the solution of the BVP consists of (10) with the  $A_{mn}$  defined in (12).

The series (11) with coefficients (12) is called a **sine series in two variables** or a **double sine series**. We summarize next the **cosine series in two variables**.

The **double cosine series** of a function  $f(x, y)$  defined over a rectangular region  $0 \leq x \leq b, 0 \leq y \leq c$  is given by

$$f(x, y) = A_{00} + \sum_{m=1}^{\infty} A_{m0} \cos \frac{m\pi}{b} x + \sum_{n=1}^{\infty} A_{0n} \cos \frac{n\pi}{c} y \\ + \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} A_{mn} \cos \frac{m\pi}{b} x \cos \frac{n\pi}{c} y,$$

where

$$A_{00} = \frac{1}{bc} \int_0^c \int_0^b f(x, y) dx dy$$

$$A_{m0} = \frac{2}{bc} \int_0^c \int_0^b f(x, y) \cos \frac{m\pi}{b} x dx dy$$

$$A_{0n} = \frac{2}{bc} \int_0^c \int_0^b f(x, y) \cos \frac{n\pi}{c} y dx dy$$

$$A_{mn} = \frac{4}{bc} \int_0^c \int_0^b f(x, y) \cos \frac{m\pi}{b} x \cos \frac{n\pi}{c} y dx dy.$$